# Half-flat causal structures and integrable systems 

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February 14, 2019

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## An outline

- Causal structures: definition, motivation, and history


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- Cayley structures


## Definitions

The conformal class of a pseudo-Riem metric $g$ on $M^{n+1}$ is uniquely determined by its field of null cones

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Roughly speaking, if $\mathcal{C}_{x}$ not quadratic, then $\mathcal{C}$ is a field of proj hypersurfaces locally described by

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If the projective 2 nd fund form of $\mathcal{C}_{x} \subset \mathbb{P} T_{x} M, \forall x \in M$ is non-degenerate everywhere one obtains a causal structure.

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Roughly speaking, if $\Sigma_{x}$ not quadratic one has a (local) Finsler metric

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T M \supset \Sigma^{2 n+1}=\{v \in T M \mid F(v)=1\}
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assuming radial transversality and non-deg of the 2nd fund form of $\Sigma_{x} \subset T_{x} M, \forall x \in M$.

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- The map $\pi \circ \iota: \mathcal{C} \rightarrow M$ is a submersion with connected fibers.
- In the fibration $\pi \circ \iota: \mathcal{C} \rightarrow M$, the fibers $\mathcal{C}_{x}^{n-1}:=(\pi \circ \iota)^{-1}(x)$ are mapped to immersed connected tangentially non-degenerate projective hypersurfaces via $\iota_{x}: \mathcal{C}_{x} \rightarrow \mathbb{P} T_{x} M$, i.e., they have non-deg projective 2 nd fund form everywhere.


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\begin{aligned}
& (M, \mathcal{C}) \stackrel{\text { locally }}{\cong}(\tilde{M}, \tilde{\mathcal{C}}) \\
& \text { at } x \in M, \tilde{x} \in \tilde{M}
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- Note that $\mathcal{C}$ can be open and be immersed as an open hypersurface in $\mathbb{P} T M$.
- For the local aspects of causal geometry $\iota$ can be assumed to be an embedding in a sufficiently small neighborhood of $\mathcal{C}$.


## Definitions and examples

Locally a causal structure can be expressed as

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L(x ; \lambda y)=\lambda^{r} L(x ; y) \text { for some } r \\
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Example : $L(x ; y)=\frac{1}{3}\left(y^{2}\right)^{3}+y^{0} y^{3} y^{3}-y^{1} y^{2} y^{3}$ :
Null cones are projectively equivalent to Cayley's cubic surface.

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Theorem (Hwang, 2013) Causal structures arising from smooth VMRTs are $V$-isotrivially flat.

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3-dimensional causal structures $\left(M^{3}, \mathcal{C}^{4}\right)$

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J^{1}(\mathbb{R}, \mathbb{R}) \cong \mathcal{K}^{3} \stackrel{\rho}{ } J^{2}(\mathbb{R}, \mathbb{R}) \cong \mathcal{C}^{4} \xrightarrow{\mu} M^{3} \cong \text { Space of solutions }
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This picture can be generalized to higher dimensions.

## The equivalence problem: causal geometry

At $(x ;[y]) \in \mathcal{C}$, with $\mu: \mathcal{C}^{2 n} \rightarrow M^{n+1}, \mu^{-1}(x)=\mathcal{C}_{x}^{n-1}$ define

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being a coframe on $\mathcal{C}$.

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| Hilbert form $\eta^{0}=\frac{\partial F}{\partial y^{i}} \mathrm{~d} x^{i}$ | Pojective Hilbert form $\omega^{0}=\frac{\partial L}{\partial y} \mathrm{~d} x^{i}$ |

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| Hilbert form $\eta^{0}=\frac{\partial F}{\partial y^{i}} \mathrm{~d} x^{i}$ | Pojective Hilbert form $\omega^{0}=\frac{\partial L}{\partial y^{2}} \mathrm{~d} x^{i}$ |
| $\eta^{0}:$ contact form on $\Sigma^{2 n+1}$ | $\omega^{0}:$ quasi-contact form on $\mathcal{C}^{2 n}$ |

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Loc. expressed as $F=1 \quad$ Loc. expressed as $L=0$
Hilbert form $\eta^{0}=\frac{\partial F}{\partial y^{i}} \mathrm{~d} x^{i} \quad$ Pojective Hilbert form $\omega^{0}=\frac{\partial L}{\partial y^{i}} \mathrm{~d} x^{i}$

| $\eta^{0}:$ contact form on $\Sigma^{2 n+1}$ | $\omega^{0}:$ quasi-contact form on $\mathcal{C}^{2 n}$ |
| :---: | :---: |
| $\mathrm{~d} \eta^{0}=-\zeta_{1} \wedge \eta^{1}-\cdots-\zeta_{n} \wedge \eta^{n}$, | $\mathrm{d} \omega^{0}=-\theta_{1} \wedge \omega^{1}-\cdots-\theta_{n-1} \wedge \omega^{n-1}$ |
| $\eta^{0} \wedge\left(\mathrm{~d} \eta^{0}\right)^{n} \neq 0$ | $-2 \phi_{0} \wedge \omega^{0}, \quad \omega^{0} \wedge\left(\mathrm{~d} \omega^{0}\right)^{n-1} \neq 0$ |

## Causal vs. Finsler

## Finsler

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Geodesics: integral curves of the Reeb vector field

$$
\eta^{0}(\mathbf{u})=1, \mathrm{~d} \eta^{0}(\mathbf{u}, .)=0
$$

Null geodesics: integral curves of the characteristic line field

$$
\omega^{0}(\mathbf{v})=0, \mathrm{~d} \omega^{0}(\mathbf{v}, .)=0
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## Causal vs. Finsler

## Finsler

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(Proj.) null cone bdle $\mathcal{C}^{2 n} \rightarrow M^{n+1}$ Loc. expressed as $L=0$

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\hline
\end{array}
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Geodesics: integral curves of
Null geodesics: integral curves of the Reeb vector field $\eta^{0}(\mathbf{u})=1, \mathrm{~d} \eta^{0}(\mathbf{u},)=$. the characteristic line field

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$\Sigma_{x} \subset T_{x} M$ is Legendrian $\Sigma_{x}^{n}=\operatorname{Ker}\left\{\eta^{i}\right\}$
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is well-def on $\Sigma$ (osc. quadric)
is well-def on $\mathcal{C}$ (osc. quadric)

## Causal vs. Finsler

| Finsler | Causal |
| :---: | :---: |
| Cartan's conn on $\Sigma$ | reg. norm. Cartan conn on $\mathcal{C}$ |
|  | Parabolic geometry of type |
|  | $\left(B_{n-1}, P_{12}\right),\left(D_{n}, P_{12}\right), n \geq 4$ |
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| str on $\mathcal{K}$ (space of geod) | str on $\mathcal{K}$ (space of null geod) |

Half-flatnesss in 4D conformal geometry
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and half-flatness or self-duality is defined as $W^{\text {asd }}=0$.

## Half-flatnesss in 4D causal geometry

Let us look at the structure equations

$$
\begin{aligned}
\mathrm{d} \omega^{0} & =\psi_{0} \wedge \omega^{0}-\theta^{1} \wedge \omega^{2}-\theta^{2} \wedge \omega^{1}, \\
\mathrm{~d} \omega^{1} & =-\gamma^{1} \wedge \omega^{0}-\psi_{1} \wedge \omega^{1}-\theta^{1} \wedge \omega^{3}+F_{2} \theta^{2} \wedge \omega^{2}+F_{1} \theta^{2} \wedge \omega^{0} \\
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Conditions that guarantees a foliation of $\mathcal{C}$ by 3 -folds are $E_{2}=V_{4}=0$. The condition $E_{2} F_{2}=0$ implies the null cones are ruled. If $E_{2}=F_{2}=0$, then $W_{4}$ and $V_{4}$ generate $W^{\text {sd }}$ and $W^{\text {asd }}$.

## Double fibrations

For 4D indefinite self-dual causal structure:

$$
\begin{aligned}
& T^{3} \hookrightarrow E_{2}, V_{4}=0 \quad \mathcal{C}^{6} \xrightarrow[E_{2}, F_{2}=0]{ } M^{4} \\
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If $E_{2}=W_{4}=V_{4}=0$ then $T$ has a projective str.
Theorem
indefinite half-flat causal on $M^{4} \Longleftrightarrow$ path geom. on $T^{3}$

## Principal null planes

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Taking d:

$$
\begin{gathered}
\mathrm{d} \lambda \equiv \lambda^{2} \theta^{1}+\lambda \psi_{2}-\gamma^{2} \bmod \omega^{0}, \omega^{1} \\
d^{2}(\lambda)=0 \Rightarrow F_{2} \lambda^{2}+2 F_{1} \lambda+F_{0}=0
\end{gathered}
$$

where $\frac{\partial}{\partial \theta^{1}} F_{i}=i F_{i-1}$.

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Let $\left\{\omega^{3}-\lambda \omega^{1}=0, \omega^{2}-\lambda \omega^{0}=0\right\}$ be an integrable $\beta$-plane.
Taking d:

$$
\begin{gathered}
\mathrm{d} \lambda \equiv \lambda^{2} \theta^{1}+\lambda \psi_{2}-\gamma^{2} \bmod \omega^{0}, \omega^{1} \\
d^{2}(\lambda)=0 \Rightarrow F_{2} \lambda^{2}+2 F_{1} \lambda+F_{0}=0
\end{gathered}
$$

where $\frac{\partial}{\partial \theta^{1}} F_{i}=i F_{i-1}$. The condition $d\left(F_{2} \lambda^{2}+2 F_{1} \lambda+F_{3}\right)=0$ implies

$$
\begin{gathered}
W_{4} \lambda^{4}+4 W_{3} \lambda^{3}+6 W_{2} \lambda^{2}+4 W_{1} \lambda+W_{0}=0, \quad \frac{\partial}{\partial \theta^{\mathrm{I}}} W_{i}=i W_{i-1} \\
f_{3} \lambda^{3}+3 f_{2} \lambda^{2}+3 f_{1} \lambda+f_{0}=0, \quad f_{3}=\frac{\partial}{\partial \omega^{3}} F_{2}, \frac{\partial}{\partial \theta^{\mathrm{I}}} f_{i}=i f_{i-1}
\end{gathered}
$$

## Principal null planes

When $E_{2}=V_{4}=0$ then $\left\{\omega^{0}, \omega^{2}, \theta^{2}\right\}$ is integrable.
The ruling planes $\left\{\omega^{0}=0, \omega^{2}=0\right\}$, are the $\alpha$-planes for the deg metric

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Thus quartic polynomial in conformal geometry is replaced by three polynomials in the causal setting

## Submaximal indefinite 4D causal str

Theorem : The submaximal 4D causal str of indefinite signature that does not descend to a conformal structure is $V$-isotrivially flat where $V$ is the Cayley cubic and its infinitesimal symmetry algebra is 8 -dimensional and solvable.

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Note that classically the Cayley cubic is associated to pair of ODEs:

$$
z_{1}^{\prime \prime}=z_{2}, \quad z_{2}^{\prime \prime}=0
$$

It appears that this pair is point equivalent to Egorov projective structure.

## Cayley structures

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The 4D structure group preserving $[\rho]$ is given by

$$
\left(\begin{array}{cccc}
\mathbf{f}_{\mathbf{0}} & \mathbf{u} & \mathbf{v} & \frac{1}{\mathbf{f}_{\mathbf{0}}} \mathbf{u v}-\frac{1}{3 \mathbf{f}_{\mathbf{0}}{ }^{2}} \mathbf{u}^{3} \\
0 & \mathbf{f}_{\mathbf{0}} \mathbf{f}_{\mathbf{1}} & \mathbf{u} \mathbf{f}_{\mathbf{1}} & \mathbf{v} \mathbf{f}_{\mathbf{1}} \\
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$$

Thus, any Cayley structure is equipped with an invariant flag given by

$$
P_{3}=\left\{\omega^{3}=0\right\} \supset P_{2}=\left\{\omega^{3}=\omega^{2}=0\right\} \supset P_{1}=\left\{\omega^{3}=\omega^{2}=\omega^{1}=0\right\}
$$

## Half-flat Cayley structures

The first order structure equations for half-flat Cayley structures is

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\mathrm{d} \omega^{0}= & -\phi_{0} \wedge \omega^{0}-\theta^{2} \wedge \omega^{1}-\theta^{1} \wedge \omega^{2} \\
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\mathrm{~d} \omega^{2}= & -\left(\phi_{0}+2 \phi_{1}\right) \wedge \omega^{2}-\theta^{2} \wedge \omega^{3}+c a_{6} \omega^{0} \wedge \omega^{1}+a_{2} \omega^{0} \wedge \omega^{3}, \\
\mathrm{~d} \omega^{3}= & -\left(\phi_{0}+3 \phi_{1}\right) \wedge \omega^{3}+a_{6} \omega^{0} \wedge \omega^{2}+a_{4} \omega^{0} \wedge \omega^{3} \\
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If $a_{1}=0$, one obtains a class of path geometries depending on 6 functions of 2 variables.
If $W_{4}=0$, one obtains a class of projective structures depending on 2 constants.

## Half-flat Cayley str: A zoo of geometric structures

Vanishing of $a_{i}$ implies integrability of $P_{2}, P_{3}$ in

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## DFK construction of a Lax pair

A torsion-free pair of ODEs $z_{i}^{\prime \prime}=F_{i}\left(t, z, z^{\prime}\right)$, defines a half-flat conformal structure on its solution space. Dunajski, Ferapontov and Kruglikov gave the following construction of a Lax pair for them

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The 4 -fold $J$ has a conformal structure with $\alpha$-planes given by

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By half-flatness the lift $\alpha$-surfaces foliate the circle bundle $N=\mathbb{P}^{1} \times M$ of $\alpha$-planes.

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The lift of $\left\{V_{1}, V_{2}\right\}$ to $N$ is ambiguous up to $\frac{\partial}{\partial \lambda}$.

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$\left[Z_{1}, Z_{2}\right] \equiv 0$ modulo $\frac{\partial}{\partial \omega^{\omega}}, \frac{\partial}{\partial \omega^{2}}$ gives two PDEs of order 3 in four variables which is involutive. The solutions depend on 6 functions of 3 variables.

Half-flat Cayley str: characterizing Fel's torsion
Given a pair $\left(z^{i}\right)^{\prime \prime}=F^{i}(t, z, z)$, let $\mathrm{D}_{t}=\partial_{t}+p^{i} \partial_{z^{i}}+F^{i} \partial_{p^{i}}$, and $\mathcal{D}=\operatorname{span}\left\{\partial_{p^{1}}, \partial_{p^{2}}\right\}$.

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Theorem: There is a one to one correspondence between 3-dimensional path geometries arising from half-flat Cayley structures and point equivalence classes of pairs of second order ODEs satisfying

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As a result $\mathbf{T}^{X}=\left(\begin{array}{ll}0 & 0 \\ 1 & 0\end{array}\right)$ in some normal frame.

## Half-flat Cayley str: characterizing Fel's torsion

Given a pair $\left(z^{i}\right)^{\prime \prime}=F^{i}\left(t, z, z^{\prime}\right)$, let $\mathrm{D}_{t}=\partial_{t}+p^{i} \partial_{z^{i}}+F^{i} \partial_{p^{i}}$, and $\mathcal{D}=\operatorname{span}\left\{\partial_{p^{1}}, \partial_{p^{2}}\right\}$.
Given $X \in \operatorname{span}\left\{\mathrm{D}_{t}\right\}$, a frame $\mathbf{V}=\left(V_{1}, V_{2}\right)$ for $\mathcal{D}$ is called normal if

$$
\begin{equation*}
\operatorname{ad}_{X}^{2} \mathbf{V}+\mathbf{T}^{X} \mathbf{V} \equiv 0 \bmod \mathrm{D}_{\mathrm{t}} \tag{2}
\end{equation*}
$$

where $\mathbf{T}^{X}$ is the torsion wrt this frame.
Theorem : There is a one to one correspondence between 3-dimensional path geometries arising from half-flat Cayley structures and point equivalence classes of pairs of second order ODEs satisfying

$$
\operatorname{rank} \mathbf{T}^{X}=1, \quad \nabla_{X} \mathbf{T}^{X}=\phi \mathbf{T}^{X}, \quad \hat{\mathbb{S}}\left(\mathbf{T}^{X}\right)=0
$$

for some function $\phi$ on $J^{1}\left(\mathbb{R}, \mathbb{R}^{2}\right)$.
As a result $\mathbf{T}^{X}=\left(\begin{array}{ll}0 & 0 \\ 1 & 0\end{array}\right)$ in some normal frame. The above ODE gives

$$
V_{1}=A_{1}+\lambda B_{1}, \quad V_{2}=A_{2}+\lambda B_{2}+\frac{\lambda^{2}}{2} A_{1}+\frac{\lambda^{3}}{6} B_{1}
$$

for some vector fields $A_{1}, A_{2}, B_{1}, B_{2}$ such that $\operatorname{ad}_{X} A_{i}=\operatorname{ad}_{X} B_{i}=0$.

## Half-flat Cayley str: A Lax pair

As in the conformal case take a slice at $t=0$.

## Half-flat Cayley str: A Lax pair

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$$
\begin{aligned}
& Y_{1}=\partial_{p^{i}}, \quad Y_{2}=\partial_{p^{2}}, \\
& Z_{1}=\partial_{z^{1}}+\frac{1}{2}\left(\partial_{p^{1}} F_{1} \partial_{p^{1}}+\partial_{p^{1}} F_{2} \partial_{p^{2}}\right), \\
& Z_{2}=\partial_{z^{2}}+\frac{1}{2}\left(\partial_{p^{2}} F_{1} \partial_{p^{1}}+\partial_{p^{2}} F_{2} \partial_{p^{2}}\right),
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\end{aligned}
$$

Find a normal frame ( $V_{1}, V_{2}$ ) for ruling lines with $V_{1} \in \operatorname{Ker} \mathbf{T}$ i.e.,

$$
\begin{gathered}
V_{1}=Y_{2}+E Y_{1}+\lambda\left(Z_{2}+E Z_{1}\right) \bmod \mathrm{D}_{t} \\
V_{2}=Y_{1}+\lambda Z_{1}+\frac{\lambda^{2}}{2}\left(Y_{2}+E Y_{1}\right)+\frac{\lambda^{3}}{6}\left(Z_{2}+E Z_{1}\right) \bmod \mathrm{D}_{t}
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\end{gathered}
$$

Following the conformal case, we find the functions $m, n$ and finally they result in 7 equations of order 3 in 3 functions $E, F^{1}, F^{2}$ which is not involutive!

## Future directions

- Integrating structure equations to give explicit examples of Cayley structures
- The moduli space of rational curves with normal bundle $\mathcal{O}(1) \oplus \mathcal{O}(2)$ in a 3 -fold is equipped with a cubic form which gives Cayley cubic when restricted to certain hyperplane.
- Examples of uniruled manifolds for which the VMRTs give a Cayley structure.
- An alternative description of Cayley structures in terms of a pair of PDEs of finite type by passing to its space of null geodesics.
- Many of our discussion extends to a larger class of half-flat $V$-isotrivial causal structures.


# Thank you for your attention! 

